

DE LA RECHERCHE À L'INDUSTRIE



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TOWARDS SIMULATING THE CELL MACHINERY AT THE ATOMIC SCALE

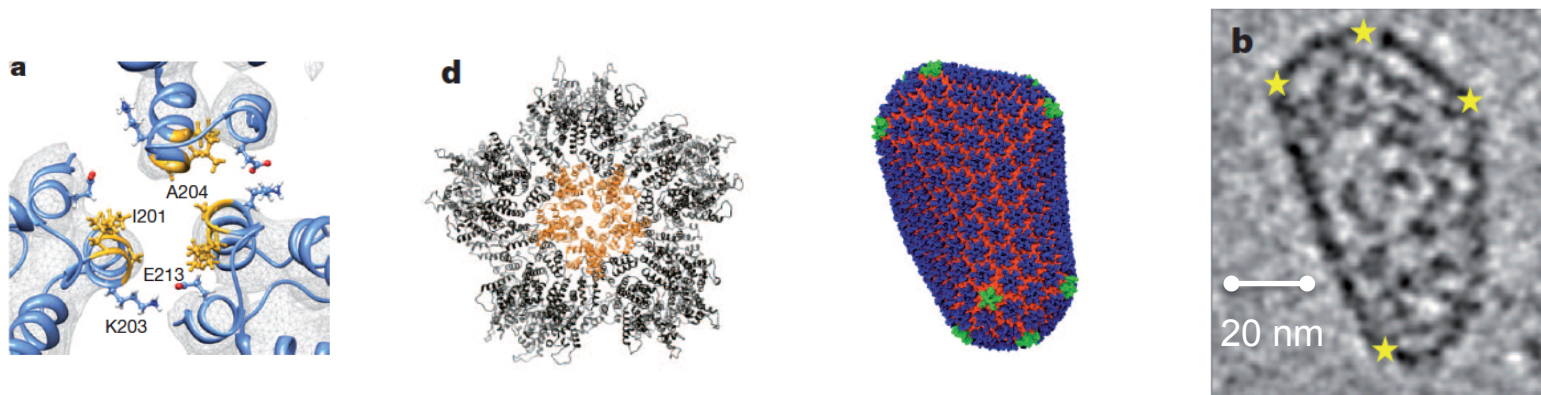
"Calculs Parallèles et Applications" | Michel Masella | Life Science Division

30 SEPTEMBRE 2014

Why to use a classical formalism and empirical potentials to model protein systems in solution: efficiency

- Much faster computations compared to quantum approaches (DFT, MP2...)
- Large systems (thousands up to millions of atoms)
- Long simulations (from the ns up to the μ s scale)

Example : simulating the mature HIV-1 capsid structure

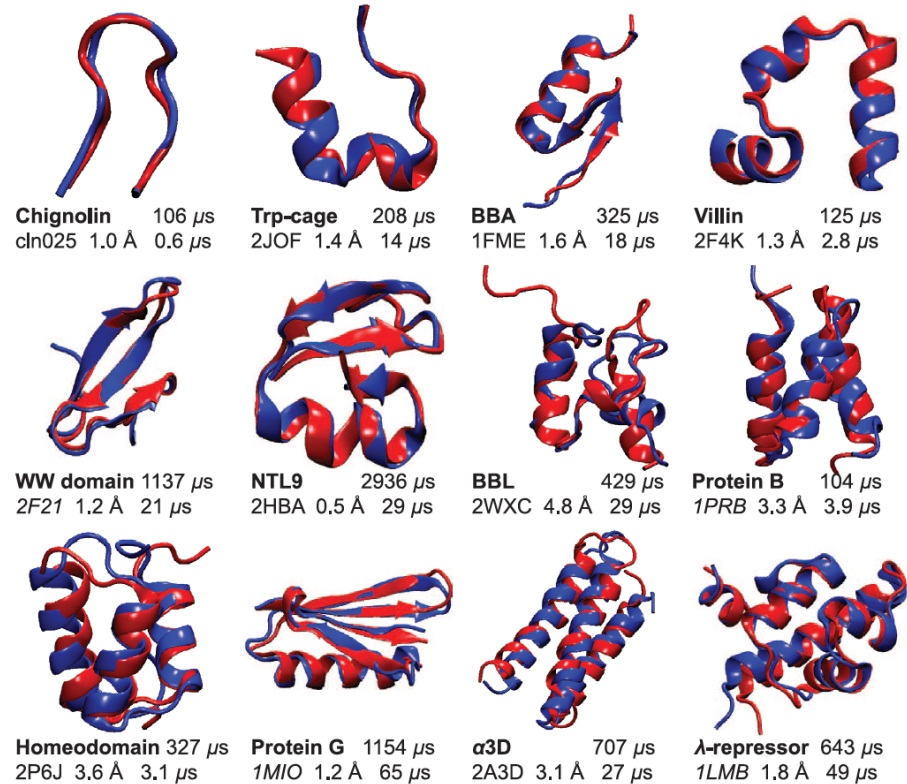


128 000 cores of the Cray “Blue Waters”, 300 ns, about $60 \cdot 10^6$ atoms ($50 \cdot 10^6 =$ water)

Zhao et al, Nature, 497 (2013) 643

Anton, a special-purpose machine for molecular dynamics simulation

Shaw et al, ACM SIGARCH, 35 (2007) 1-12



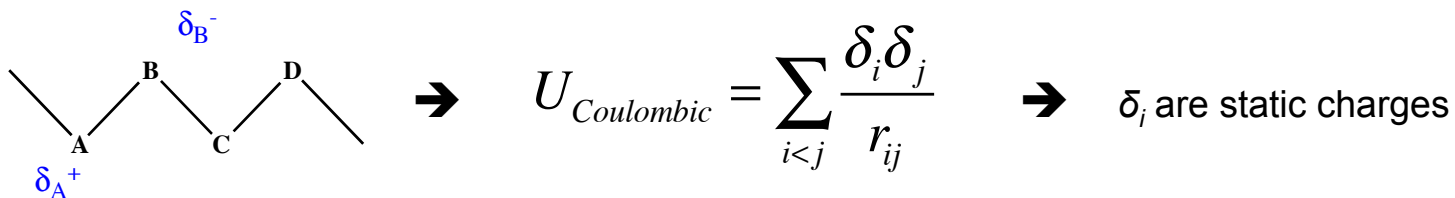
How fast folding proteins fold

Lindorff-Larssen et al, *Science*, 334 (2011) 517

The main drawback of MM approach: accuracy

Common force-fields are based on additive potentials (CHARMM, AMBER, OPLS ...)

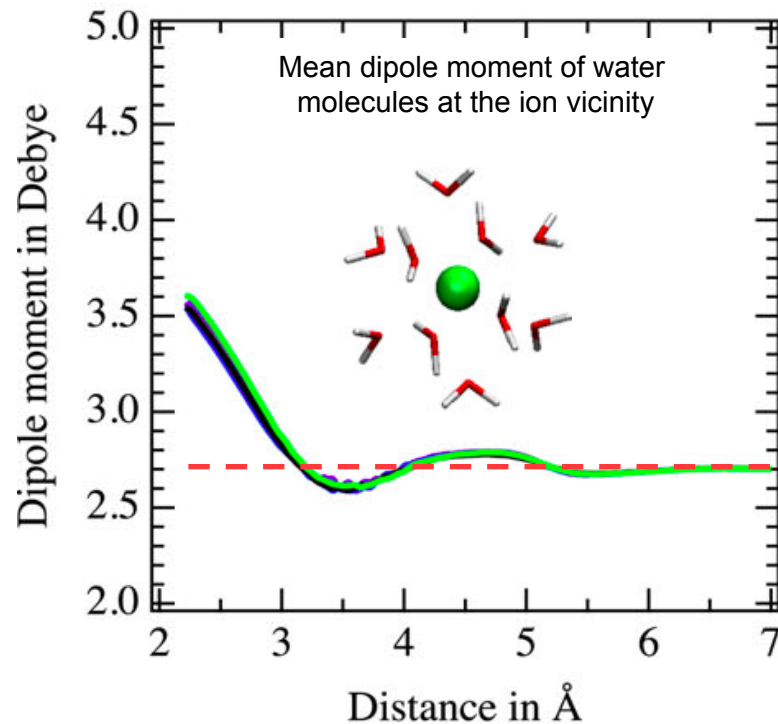
- 1 - The model parameters are constant along an MD trajectory
- 2 – Reliability of such an approach when mimicking microscopic electrostatic interactions



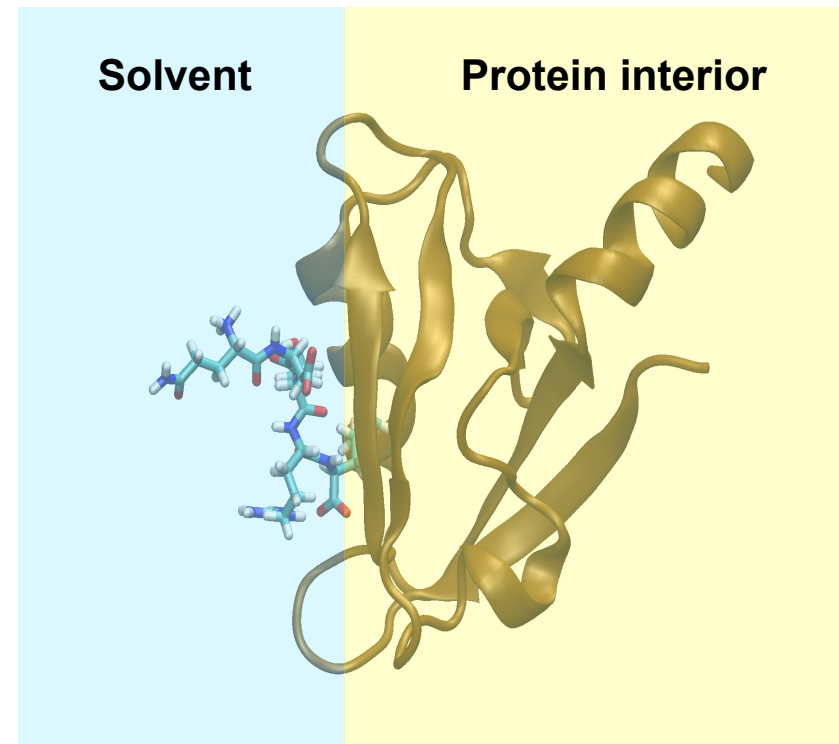
- 3 – Experimental results are commonly used for assigning/refining parameters, transferability ?
- 4 - Ok, you may say that for a homogeneous and isotropic system, you are using a mean field approach...

However, what about interfaces ?

The case of Cm(III) interacting with water



All protein/ligand interactions correspond to an interface problem



These phenomena can not be accounted for using static charges to model the molecular electronic cloud properties

They allow to account for environmental effects on electrostatic molecular properties considered in classical molecular dynamics

Many approaches have been proposed :

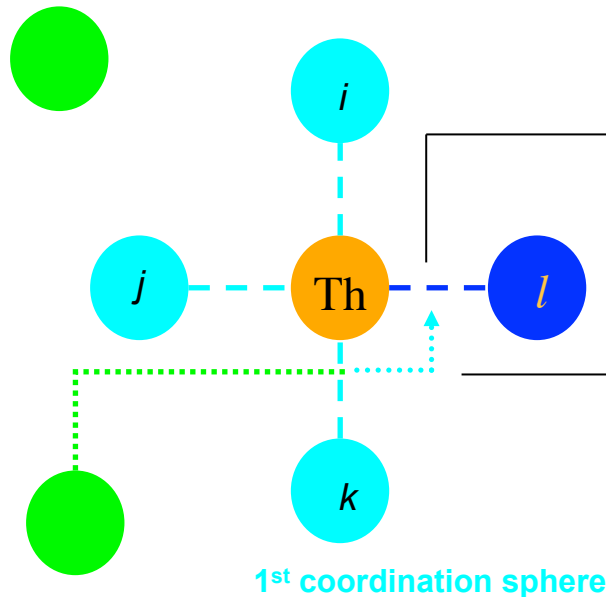
- Fluctuating charge approaches: the “static” δ_i charges are allowed to fluctuate
- Drude oscillators: “extra” atomic charges are introduced
- **Induced dipole moments: new degrees of freedom are introduced**

$$\boldsymbol{\mu}_i = \alpha_i \left(\mathbf{E}_i - \sum_{j=1, j \neq i}^N \mathbf{T} |\mathbf{r}_j - \mathbf{r}_i| \boldsymbol{\mu}_j \right) \Rightarrow \text{to be solved iteratively, } O(N^2)$$

\mathbf{T} is the dipolar tensor and α_i is the center i polarizability, usually isotropic
 \mathbf{E}_i is the electric field acting on center i
 $\boldsymbol{\mu}_j$ is the induced dipole moment on center j

Accounting for inter atomic interactions with a weak covalent character

2nd coordination sphere



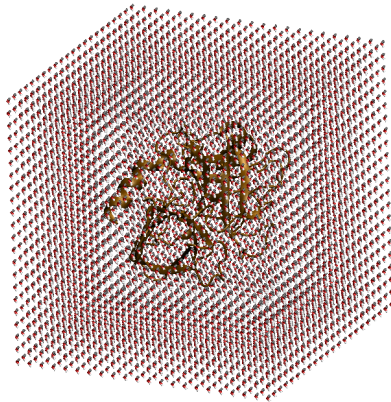
$$U_{ct} = \sum_{l \in \text{water}} D_e^c \exp(-\beta \cdot r_{Th-l})$$

$$D_e^c = D_e \left[1 - \xi \sum_{i \in \text{water}, i \neq l} \exp\left(-\frac{(r_{Th-i} - r_e)^2}{\gamma_r}\right) \right]$$

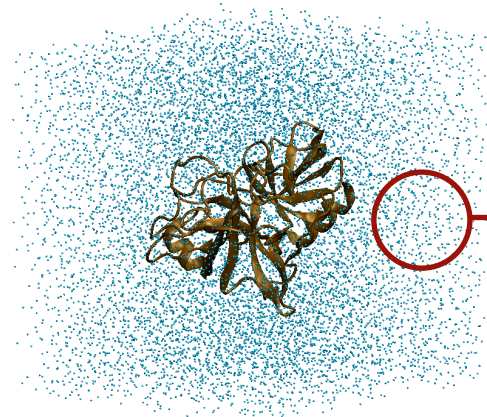
Seems computationally demanding, however scales as $O(N)$

A SOLVENT COARSE GRAINED APPROACH

Explicit solvent approach, need of FFT !



Coarse grained approach
Efficiency and microscopic solvation properties



$$\alpha_s = \frac{1}{4\pi} \left(1 - \frac{1}{\epsilon_s} \right) \times \frac{1}{\rho_s}$$

i.e. $\alpha_s \propto \Delta v_s$

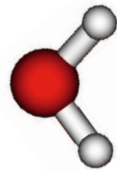
$$F_P[\mathbf{P}] = \frac{1}{2} \int \frac{\epsilon(\mathbf{r}) \mathbf{P}(\mathbf{r})^2}{\chi(\mathbf{r})} dV - \int \mathbf{P}(\mathbf{r}) \cdot \mathbf{E}_s(\mathbf{r}) dV$$

$$\mathbf{P}(\mathbf{r}) = \frac{\chi(\mathbf{r})}{\epsilon(\mathbf{r})} \mathbf{E}_s(\mathbf{r}) \longrightarrow \mathbf{p}_k^s = \Delta v \mathbf{P}(\mathbf{r}) = \frac{\mathbf{P}(\mathbf{r})}{\rho_s}$$

Implicit solvent approaches
Poisson-Boltzman
Warshell's grid approach

Haduong et al, J. Chem. Phys., 117 (2002) 541
Masella et al, J. Comput. Chem., 29 (2008) 1707

In the case of water, number of solvent atoms divided by 3



An explicit water molecule



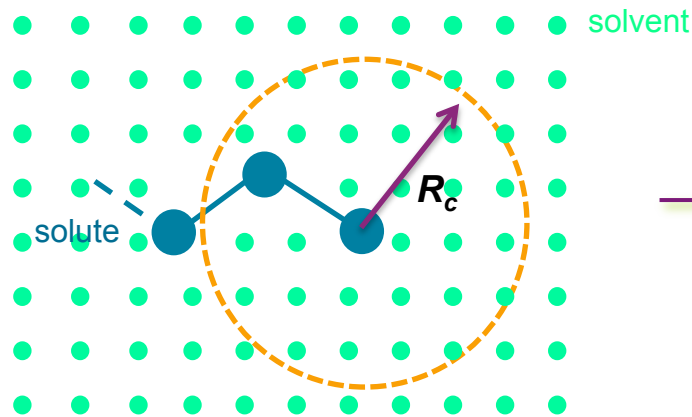
A nonpolar polarizable particle
= a water molecule



What about the liquid water hydrogen bond network ?

A solution from DPD

Systematic truncation of long tail electrostatic interactions

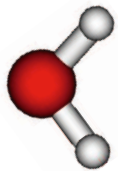


What about electrostatic long tail effects ?

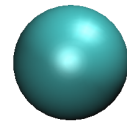
with $R_c = 12 \text{ \AA}$, the solvent cost represents 10 to 30% of the total cpu time

Particle polarizability proportional to their individual volume : $\alpha_s \propto \Delta V_s$

→ a multi-scale coarse grained approach to handle long tail bulk electrostatic



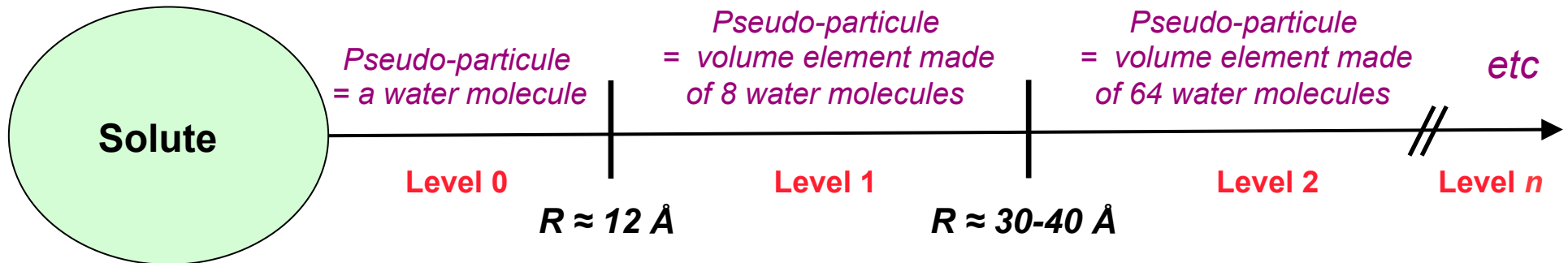
An explicit water molecule



A polarizable pseudo-particle
= a water molecule

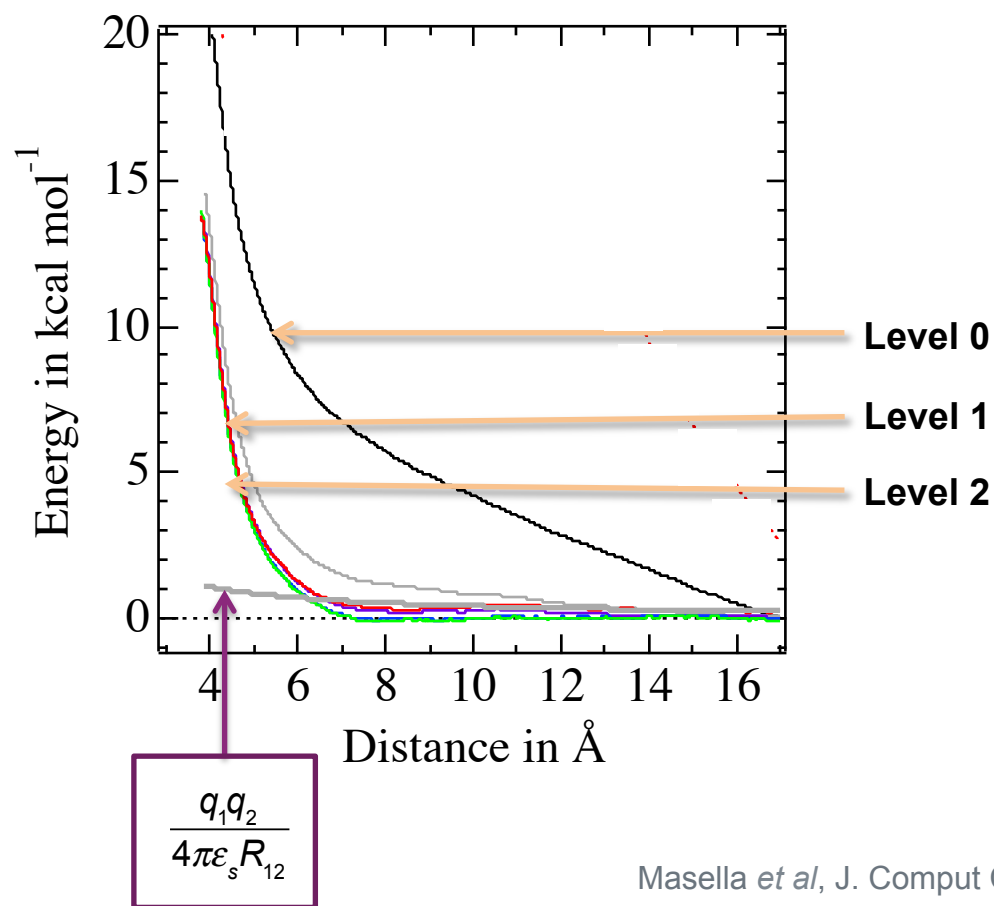


A larger polarizable pseudo-particle
= a volume element of liquid water



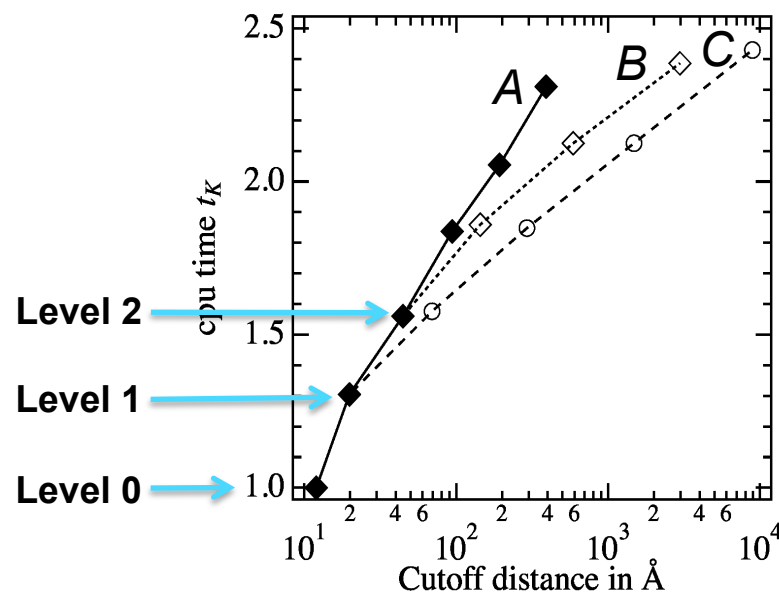
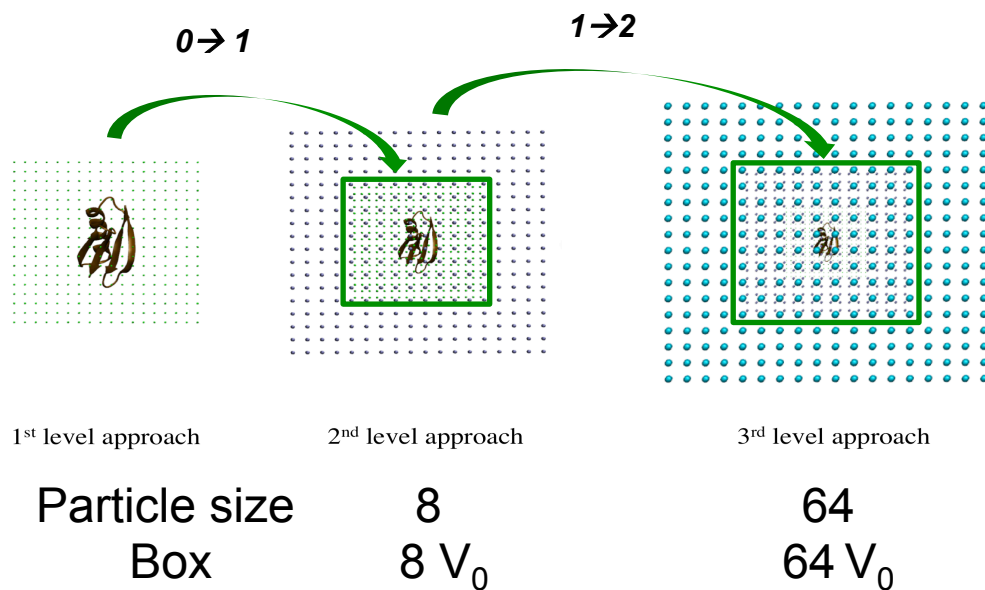
Potential of mean force of Cl-/Cl- and the coarse grained approach level

(1ns simulation, solvent boxes made of 4096 particles, umbrella sampling, etc...)



Masella *et al*, J. Comput Chem, 2013

M. Masella et al, *J. Comput Chem*, **32** (2011) 2664; *ibid*, **34** (2013) 1112

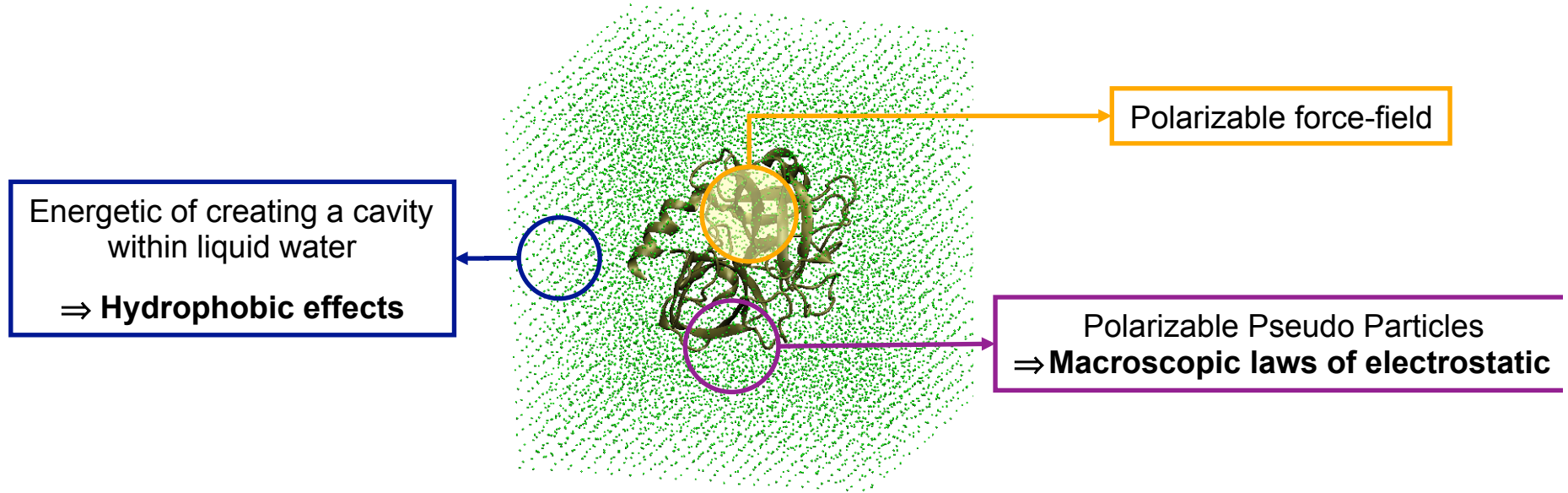


Boxes	0 → 1	1 → 2	2 → 3	3 → 4
A	x2	x2	x2	x2
B	x2	x2	x3	x4
C	x2	x3	x4	x5

Now, the bottleneck is handling the interactions within the solute...

The code POLARIS(MD) © CEA/DSV

Masella, Mol. Phys., 2006 ; Masella *et al*, J. Comput Chem, 2008; *ibid*, 2011; *ibid*, 2013



Supported by the Exascale Computing Research Laboratory



energie atomique • energies alternatives



Reference : Intel Sandy-Bridge octo-cores 2.7 Ghz (CURIE)

All atom simulations+ PME summation techniques

A water box of 1000 molecules → **3.0 ns per day** (1 cpu)
4.5 ns per day (4 cpu)

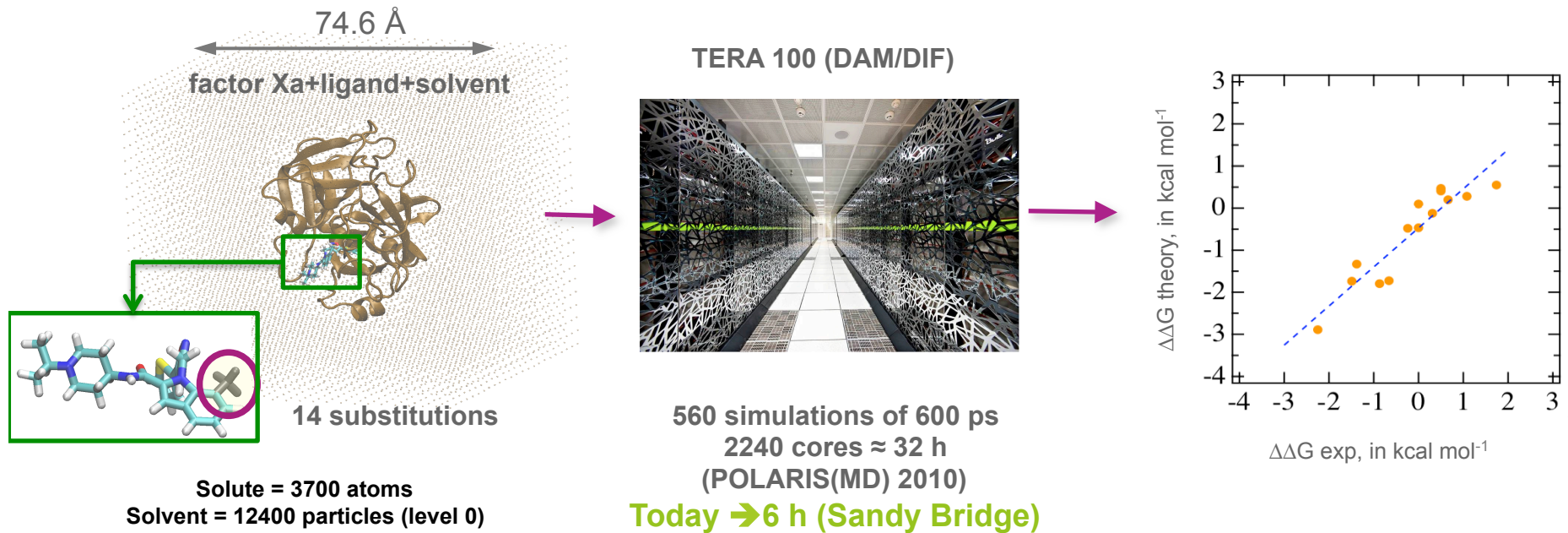
Coarse grained simulations + solute interactions $O(N^2)$

6000 particles + solute = 900 atoms → **15 ns per day** (level 0 + 1 cpu)
13 ns per day (level 1 + 1 cpu)
11.5 ns per day (level 2 + 1 cpu)
9.5 ns per day (level 3 + 1 cpu)

Largest system simulated :

8400 atoms solute + 51 000 particles, **850 ps per day** (level 0 + 1 cpu)

Free energy $\Delta\Delta G$ computations (TI schemes and 40 steps)



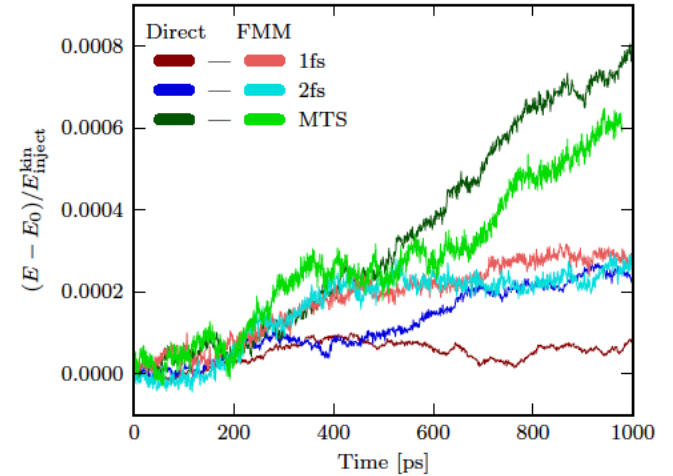
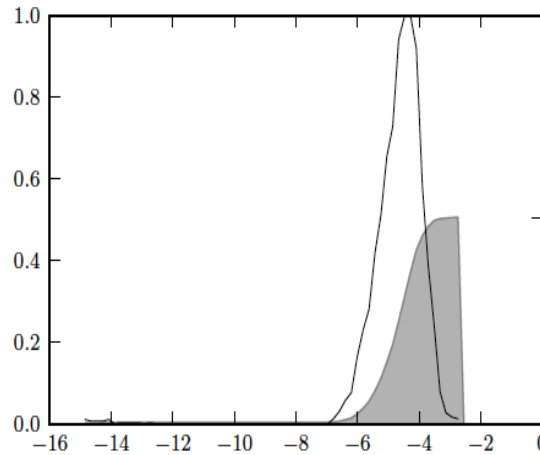
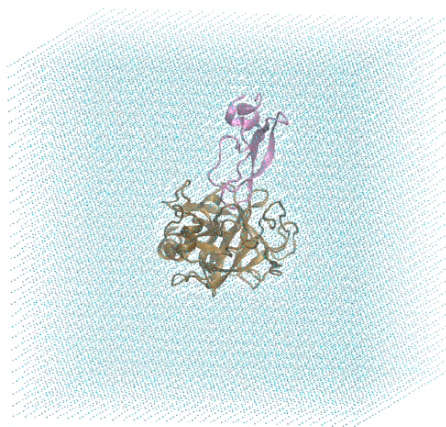
**POLARIS(MD) 2013 + solute $O(N^2)$ + 500 000 h CURIE :
 \rightarrow 400 $\Delta\Delta G$ computations**

A $O(N)$ FAST MULTIPOLE METHOD

$$\phi(\mathbf{x}_b - \mathbf{x}_a) = \frac{q_a}{|\mathbf{x}_b - \mathbf{x}_a|} \quad \longrightarrow \quad \phi(\mathbf{x}_b - \mathbf{x}_a) \approx \sum_{|\mathbf{n}| \leq p} \sum_{|\mathbf{m}| \leq p - |\mathbf{n}|} \frac{(-1)^{|\mathbf{n}|}}{n!m!} r_b^n r_a^m \nabla^{n+m} \phi(z_B - z_A)$$

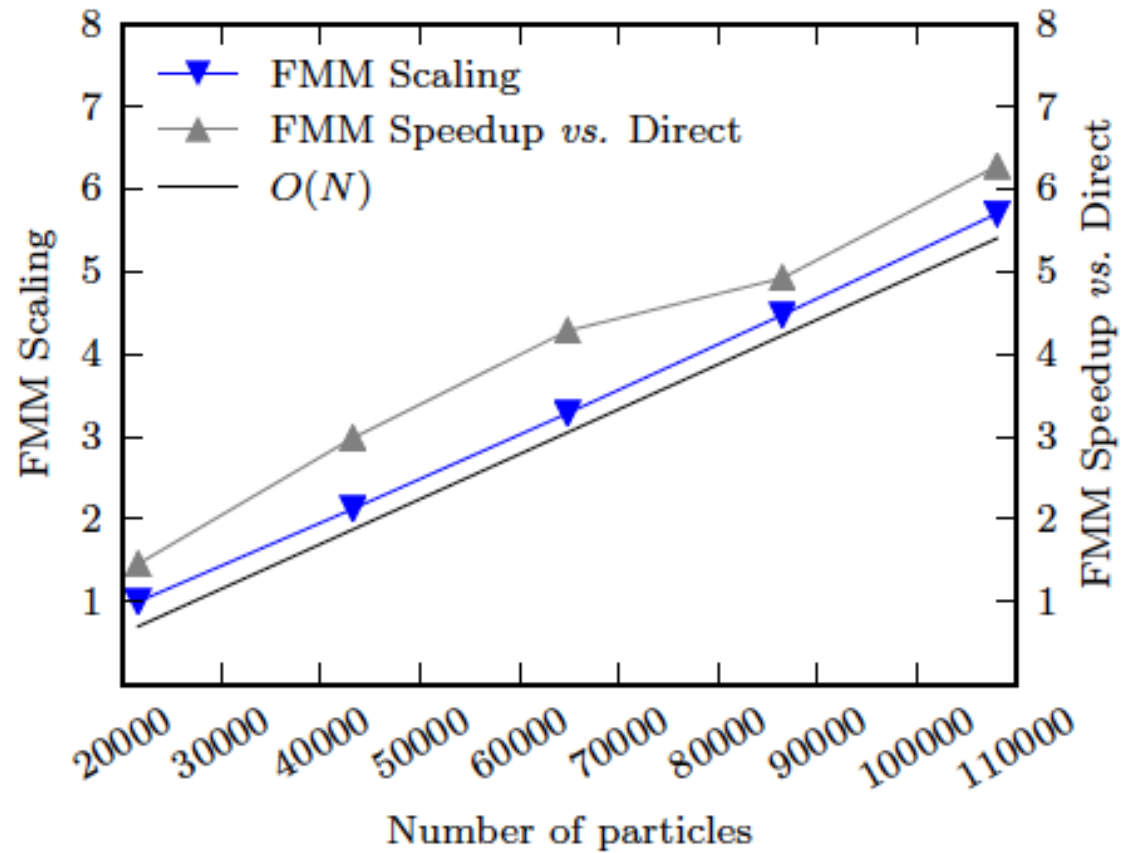
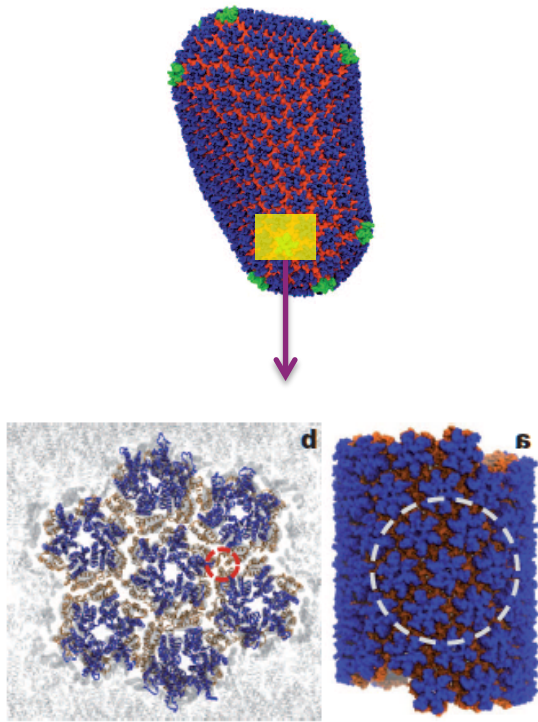
$$+ \quad \boldsymbol{\mu}_a = 2q_a^\mu \boldsymbol{\delta} \mathbf{l}_a$$

+ atoms organized via a kd-tree spatial decomposition



Coles and Masella, arXiv:1408.6527v1

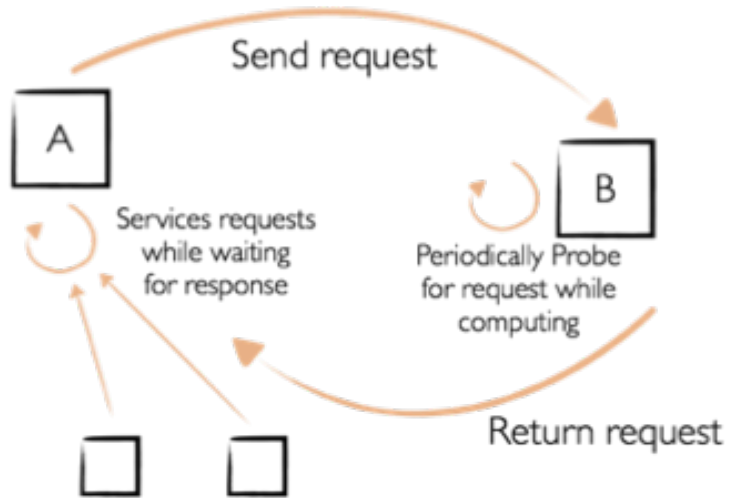
The mature HIV-1 capsid



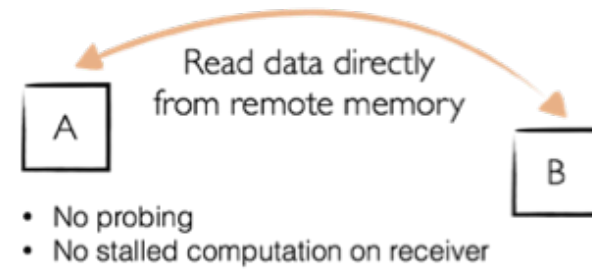
Coles and Masella, *arXiv:1408.6527v1*

TOWARDS A ONE SIDED COMMUNICATION SCHEME

Standard 2-sided communication scheme



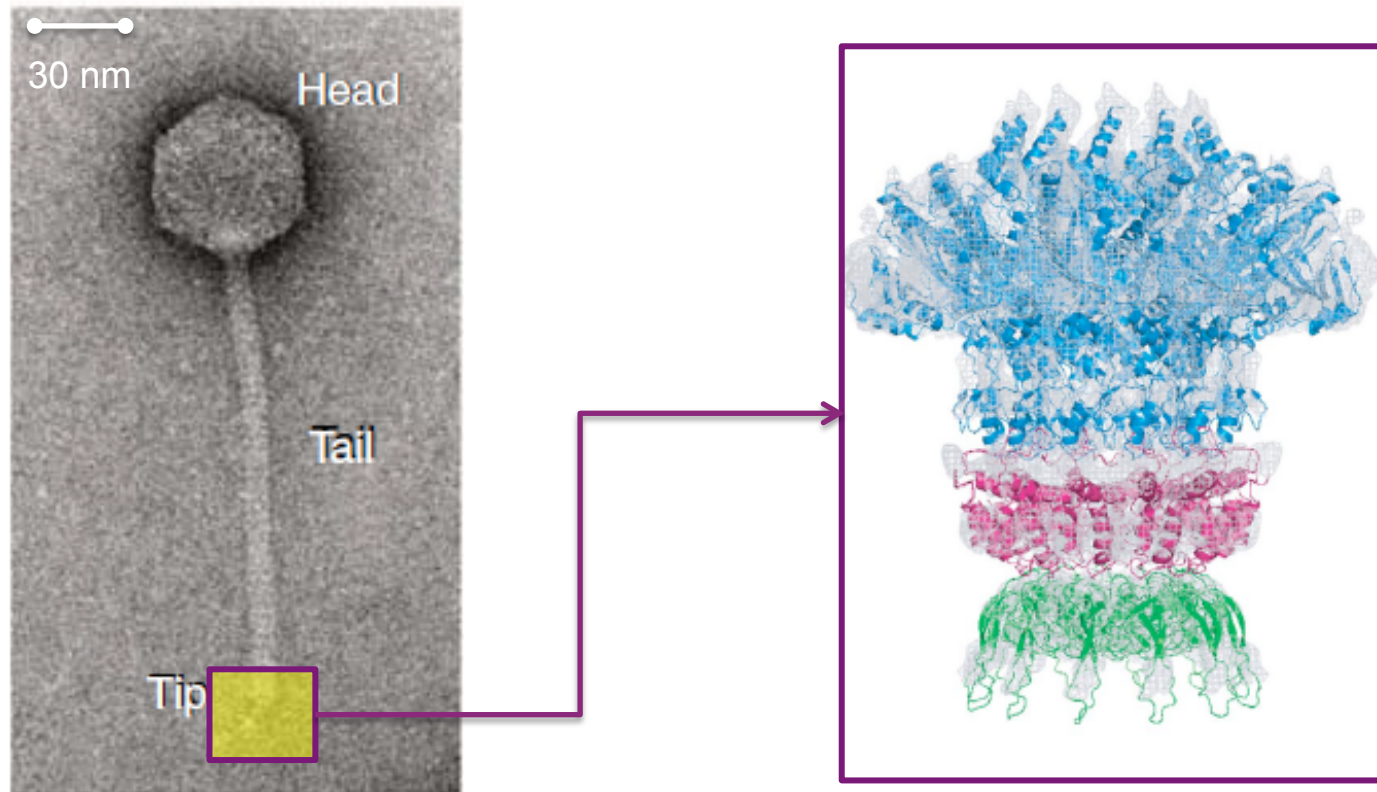
1-sided scheme, MPI v 3.0



- No probing
- No stalled computation on receiver

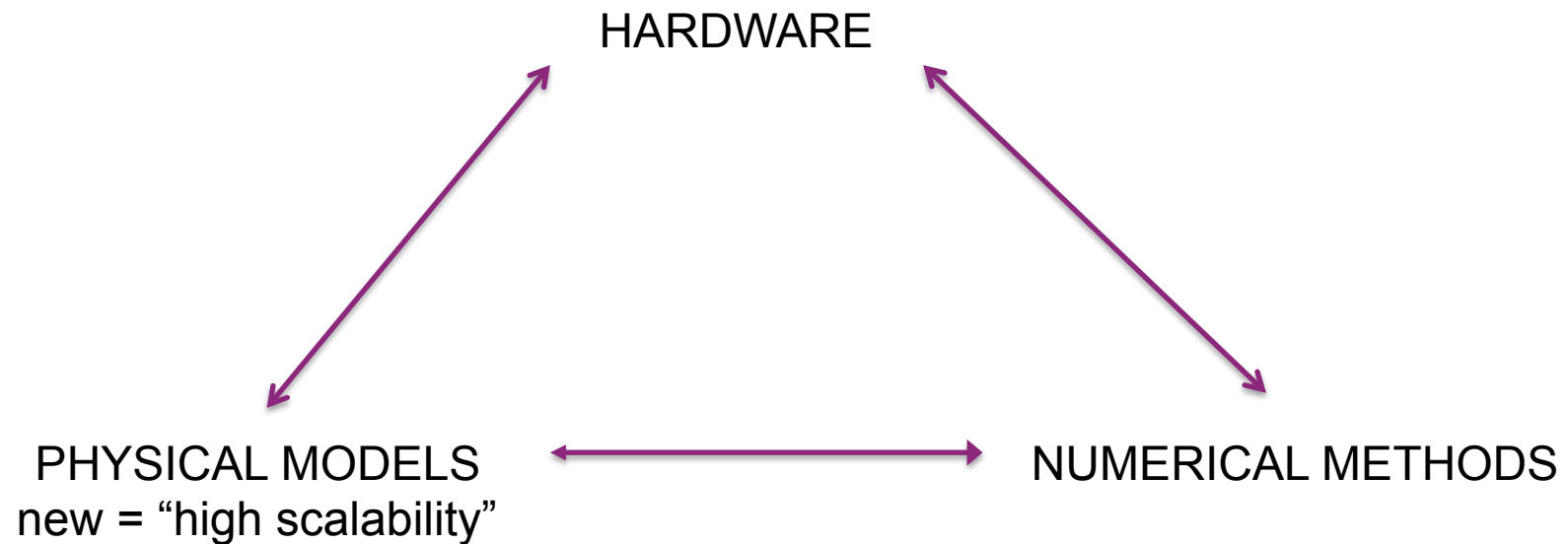


Bacillus subtilis bacteriophage SPP1 structure



Lhuillier et al, PNAS, 106 (2008) 8507

TOWARDS EXAFLOPIC SIMULATIONS



30 SEPTEMBRE 2014

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